

# Massive stars from various simulations: different, but why?

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Seminar of the Institute of Astronomy  
23 May 2022

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  - libraries / grids, e.g. Geneva models, Bonn models...

massive:  $> 8 M_{\odot}$

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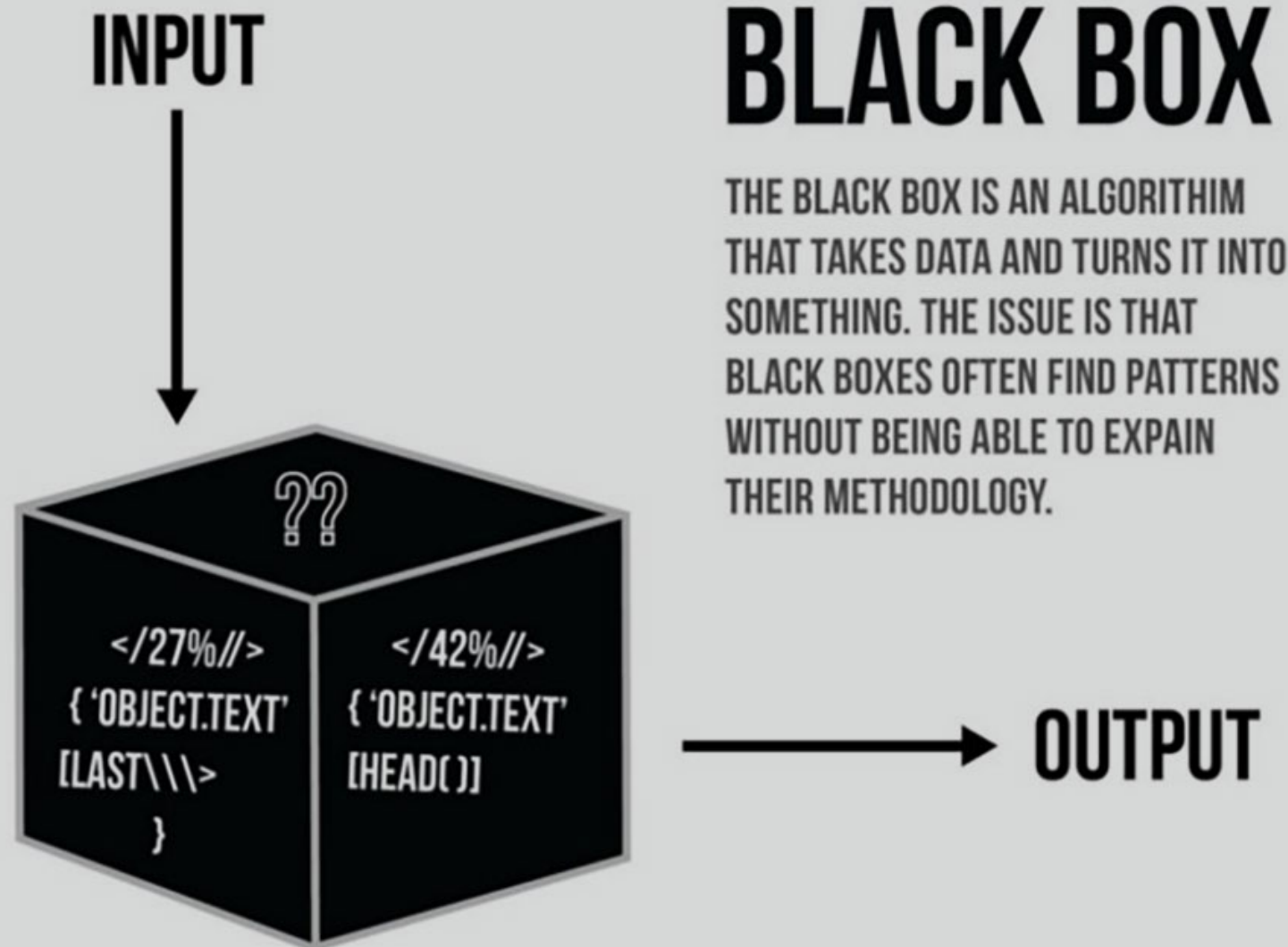
- ...maybe even you?
- *Massive* star models (“tracks”):
  - libraries / grids, e.g. Geneva models, Bonn models...
- Really wide range of usage:
  - obtaining mass & age of observed stars
  - star-formation simulations, starcluster formation studies
  - chemical evolution of the Universe
  - binary population synthesis → gravitational-wave event rates

massive:  $> 8 M_{\odot}$

just  
examples,  
there are  
more

What do  
you do?

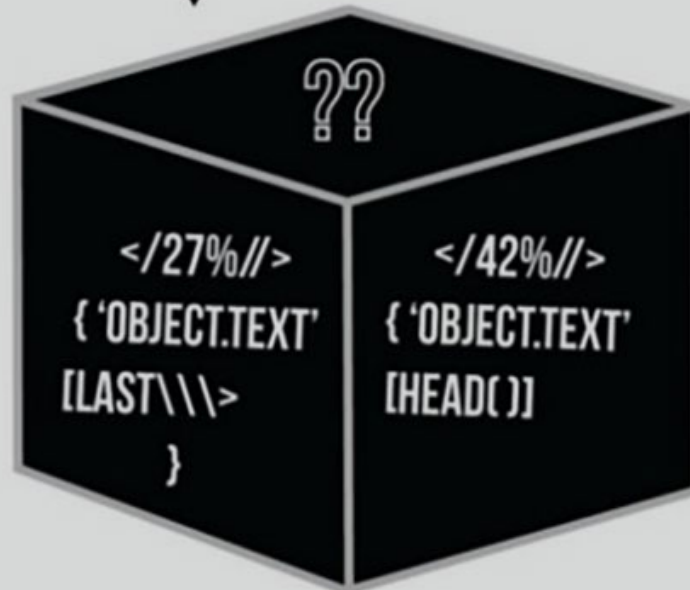
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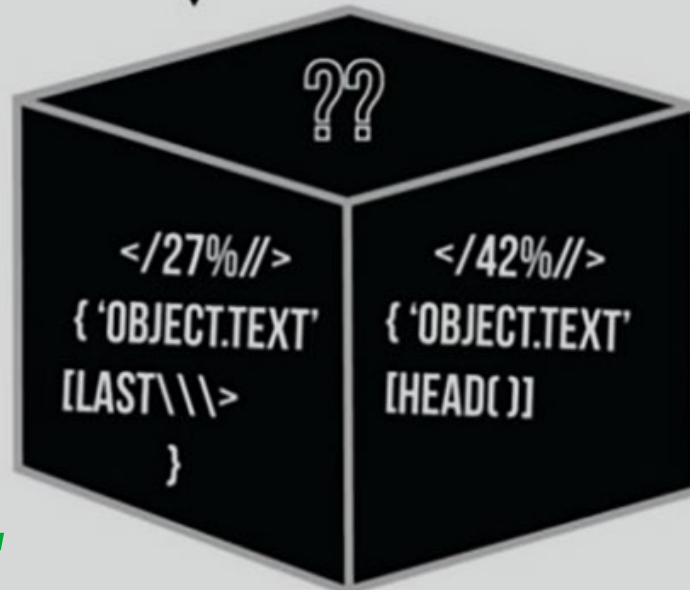
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OUTPUT

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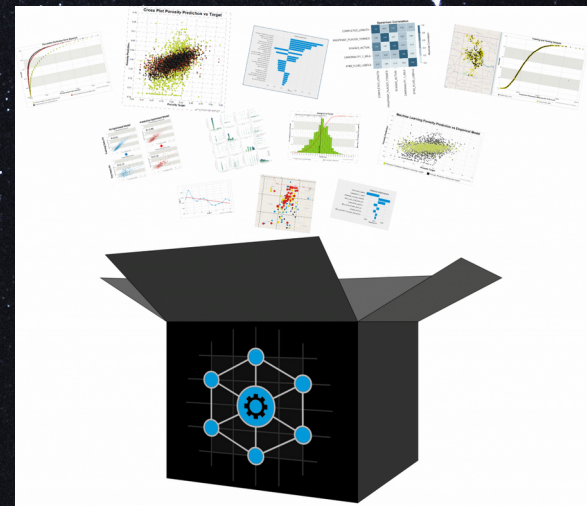
*However...*



*Let's peek into to box!*

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Agrawal & Szécsi et al. (2022, MNRAS)



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We compare 5 sets of stellar evolutionary models from 5 independent projects

– so that you don't have to ;)

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- PARSEC (Padova code)
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*Only comparing:*  
**models with the same mass and composition\***  
(single stars with no or slow rotational rate)

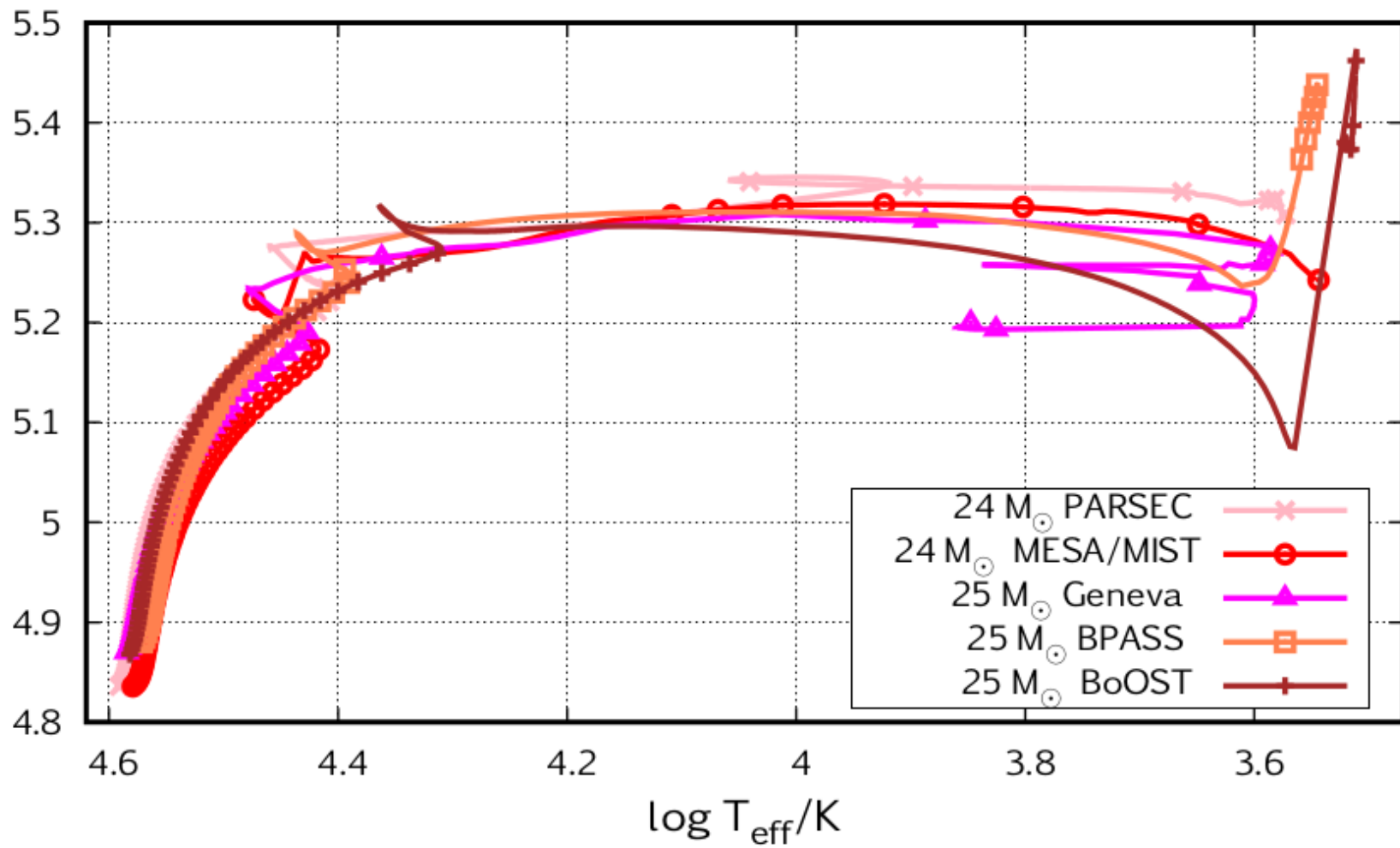
\*namely, Solar

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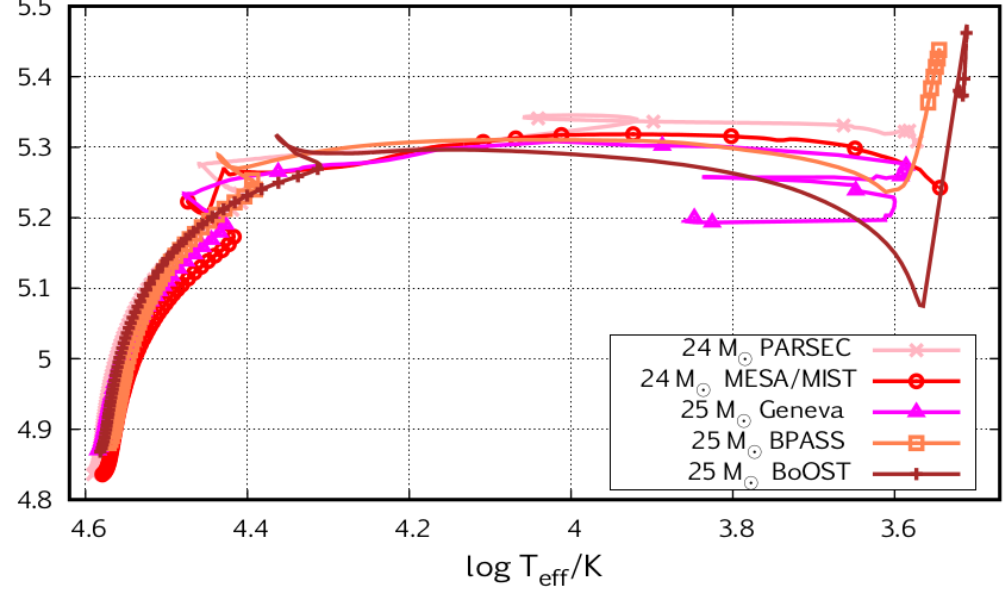
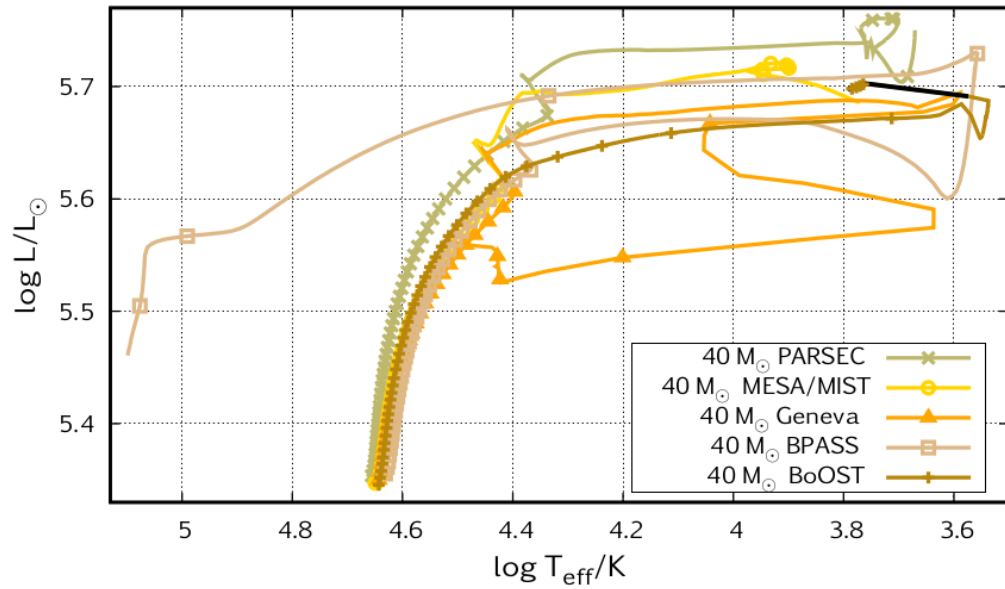
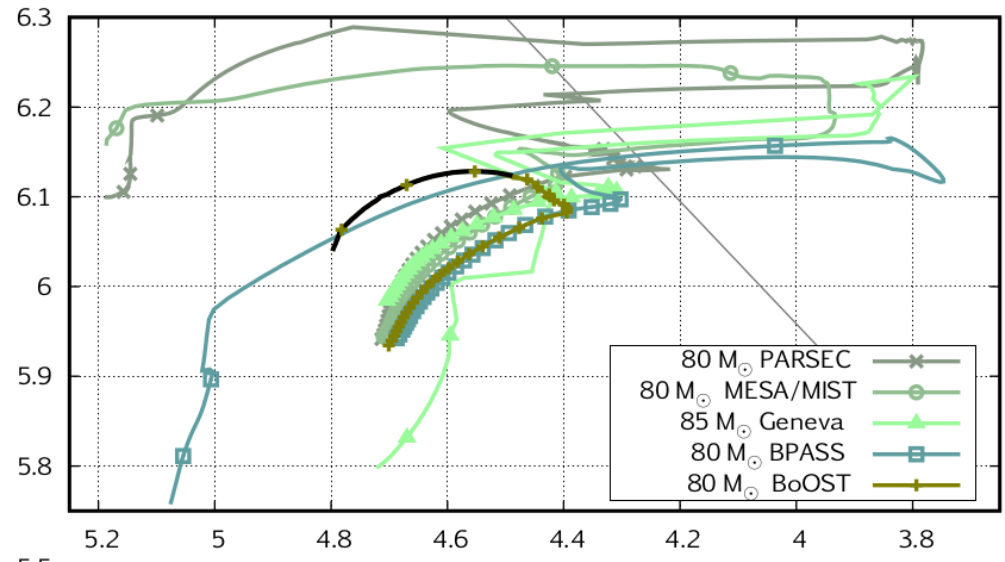
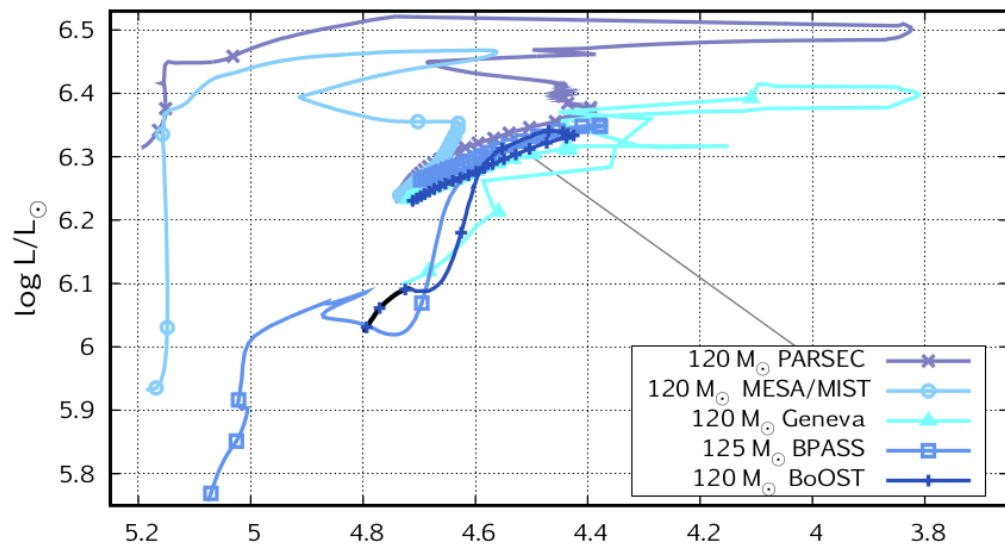
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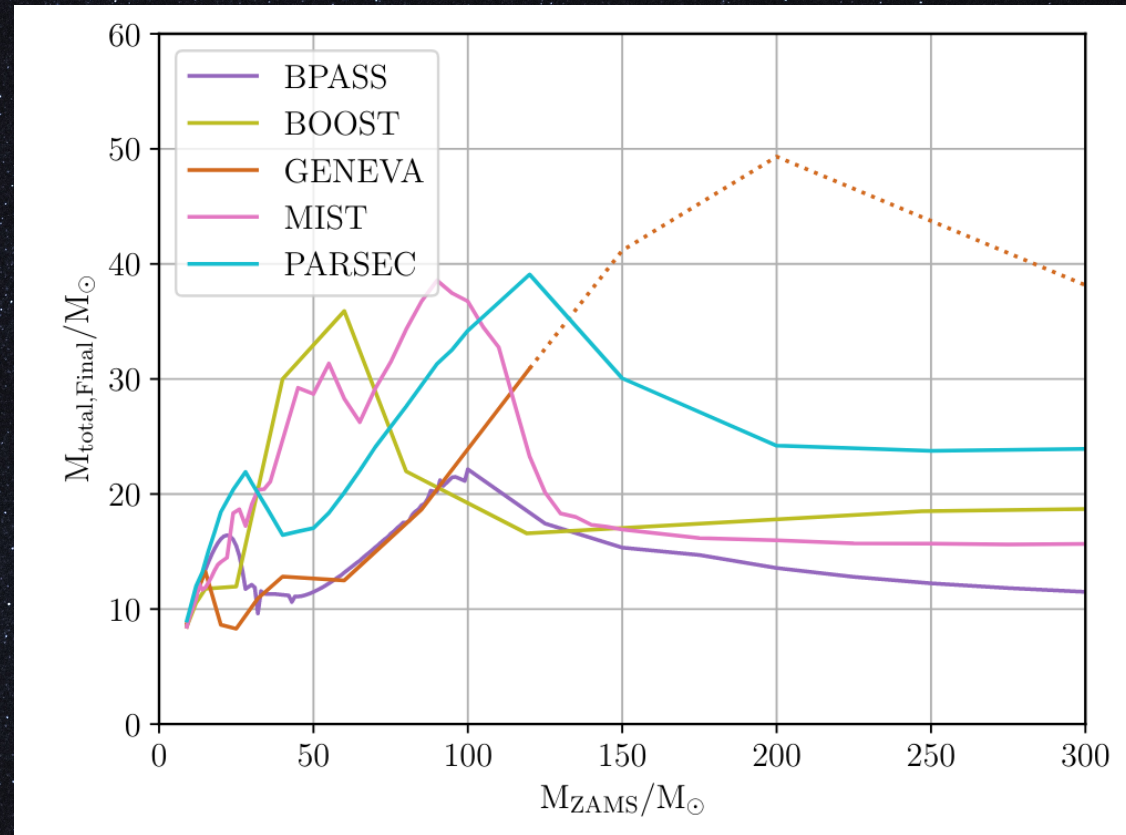
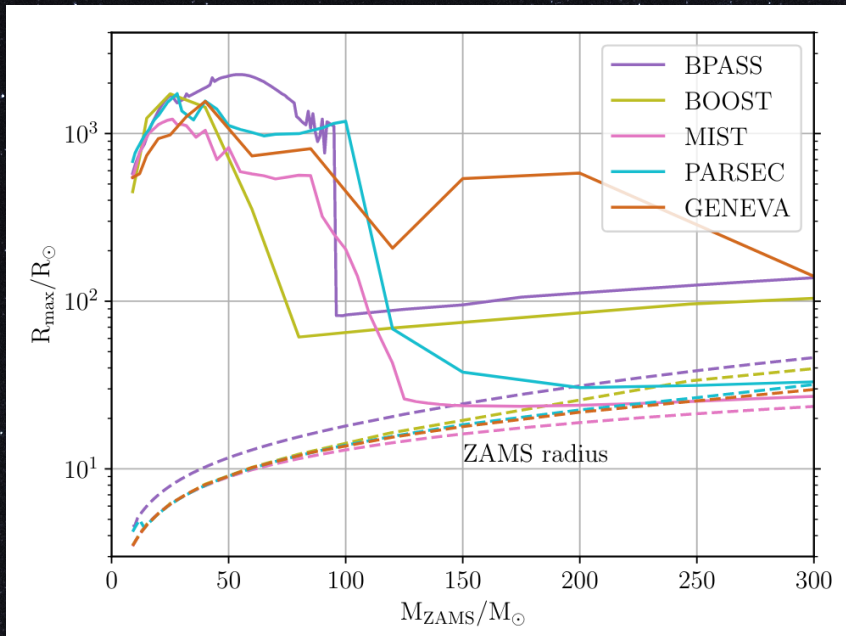
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# *What about other predictions?*

P. Agrawal (2021, *PhD thesis*)

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# What about other predictions?



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*O-okay, but... why??*

P. Agrawal (2021, *PhD thesis*)

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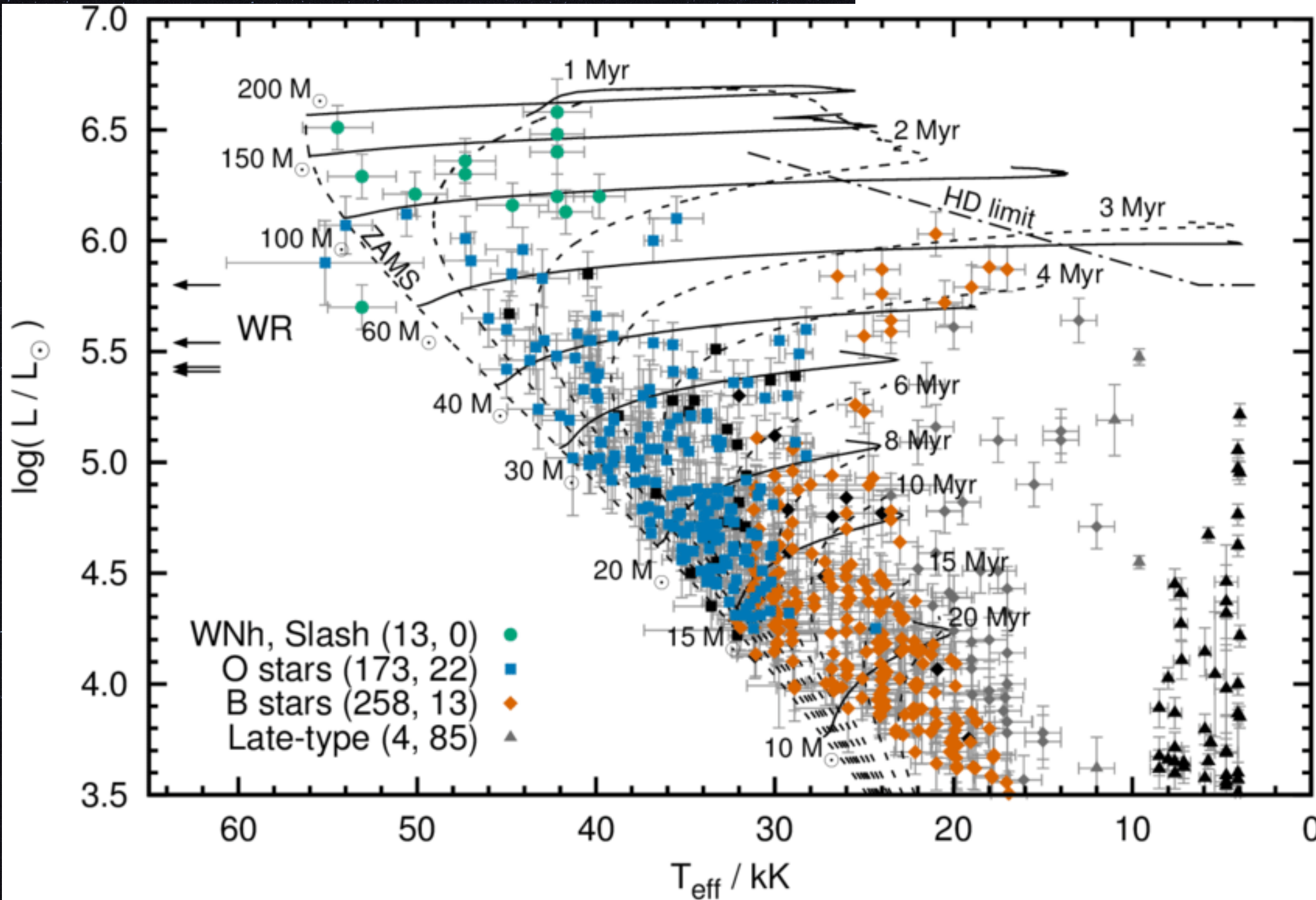
*Quick and dirty answer:*

**we don't really  
understand  
massive star physics  
that well. (Yet.)**

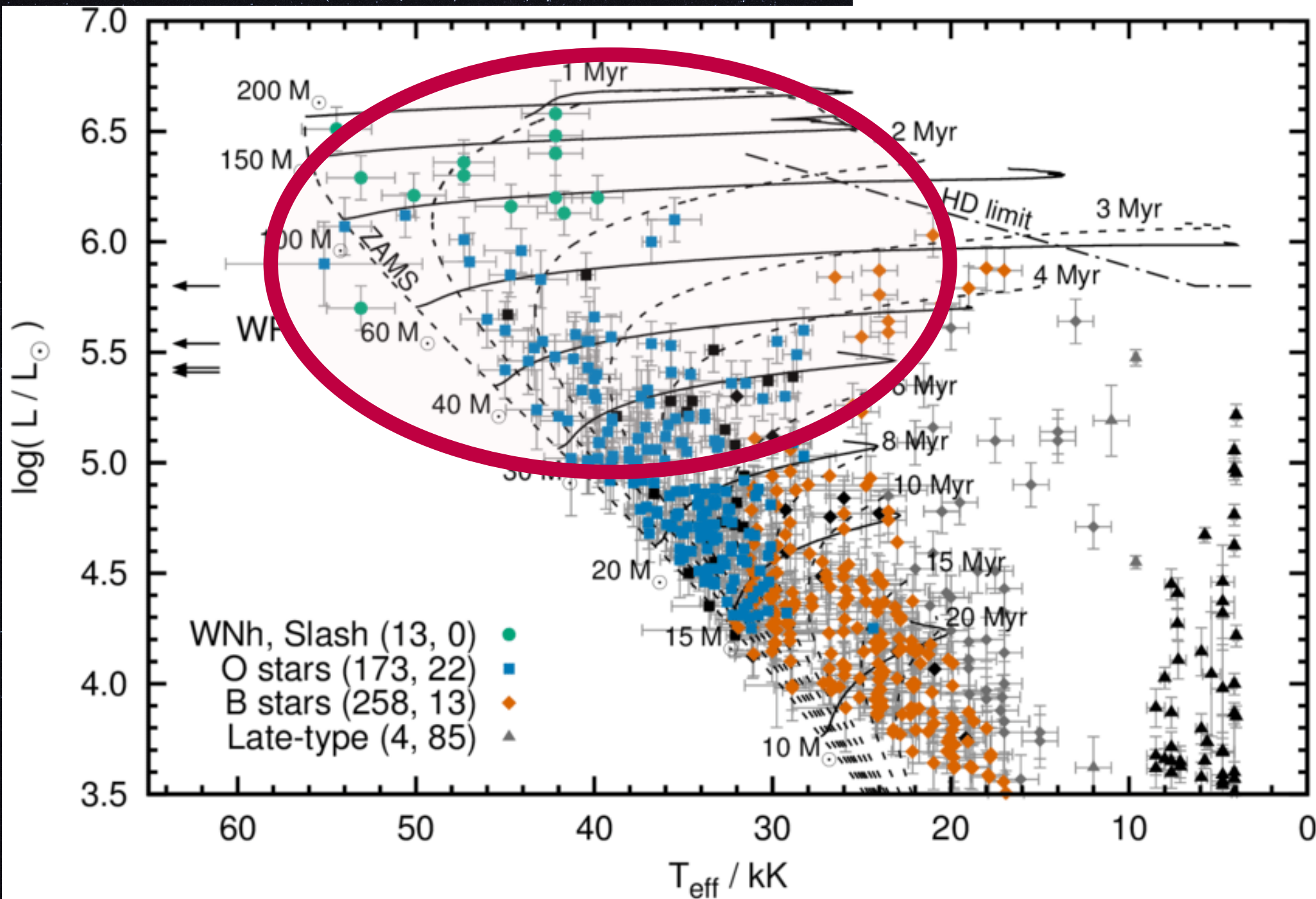
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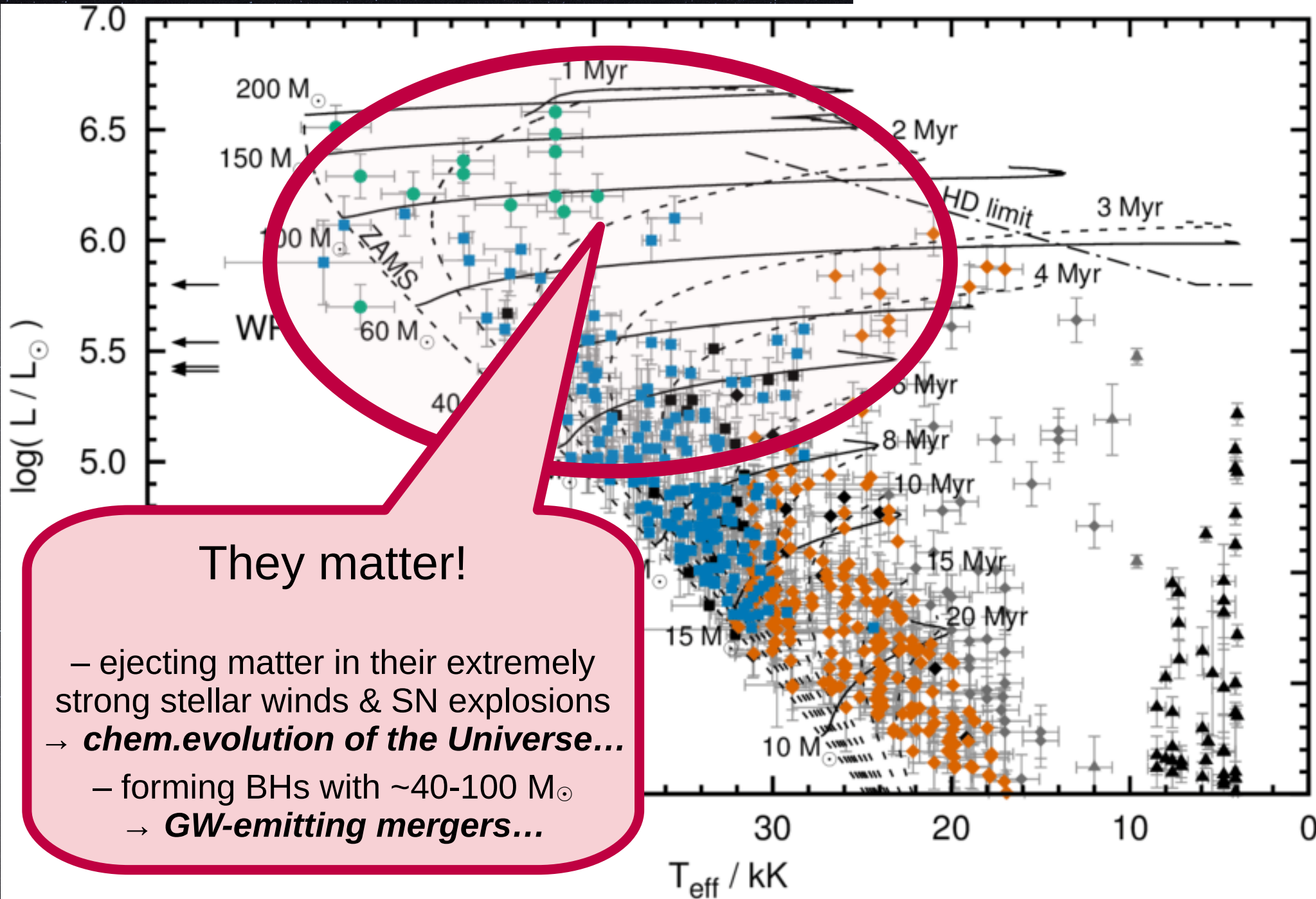
30 Doradus star-cluster in the Large Magellanic Cloud galaxy  
(VFTS survey, 2018)



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They matter!

- ejecting matter in their extremely strong stellar winds & SN explosions  
→ **chem.evolution of the Universe...**
- forming BHs with  $\sim 40\text{-}100 M_{\odot}$   
→ **GW-emitting mergers...**

***Again...***  
***different, but why??***

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***Again...  
different, but why??***

***Long answer...***

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When the equilibrium\* is compromised:

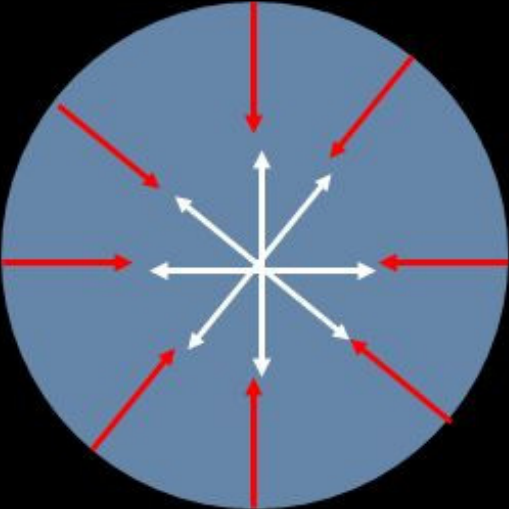
## the Eddington limit

\* *between  
gravity & radiation pressure*

# Eddington limit

Radiative Force

Gravitational Force

$$g_{rad} = \int_0^{\infty} d\nu \frac{\kappa_{\nu} F_{\nu}}{c}$$


$\frac{GM}{r^2}$

$$\Gamma_e \equiv \frac{g_e}{g} = \frac{\kappa_e L / 4\pi r^2 c}{GM / r^2} = \frac{\kappa_e L}{4\pi GMc}$$

Credit: Stan Owocki

# *Other reasons for falling out of equilibrium:*

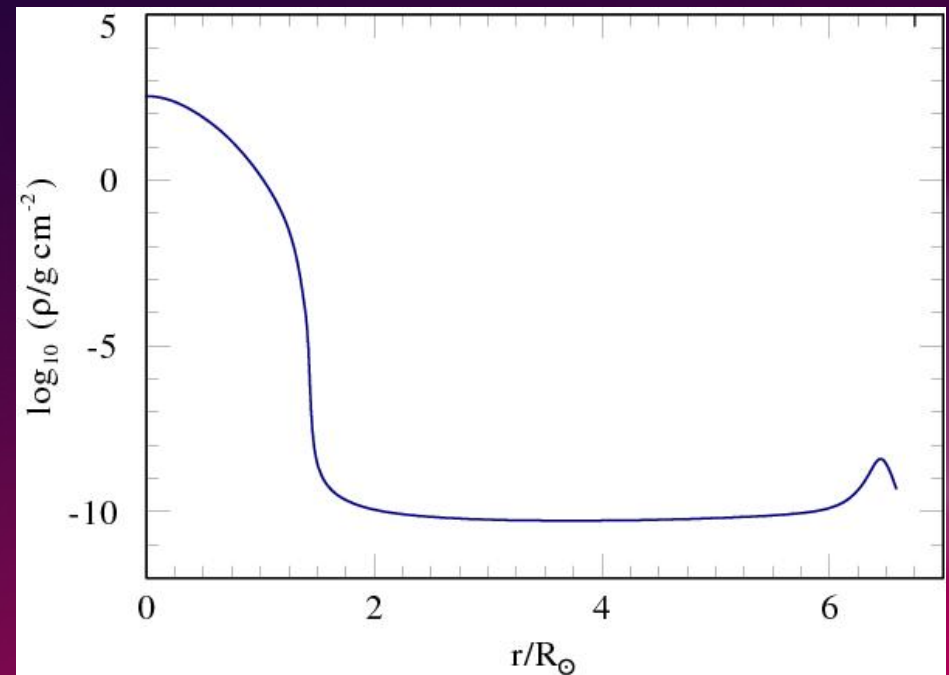
- iron core
  - gravitational collapse & SN (due to bounce-back)
- pair-instability
  - grav. collapse & subsequent thermonuclear explosion (PISN) or pulsations (puls-PISN)
- end of a burning phase
  - restructuring, crossing the Herzprung-gap...
- ...

of approaching the Eddington-limit

# Consequences for the stellar interior

- density (and pressure) inversion *in the envelope*
- no efficient energy transport mechanism here (weak convection)
- → envelope “inflation”
- numerical difficulties...

*density inversion:*



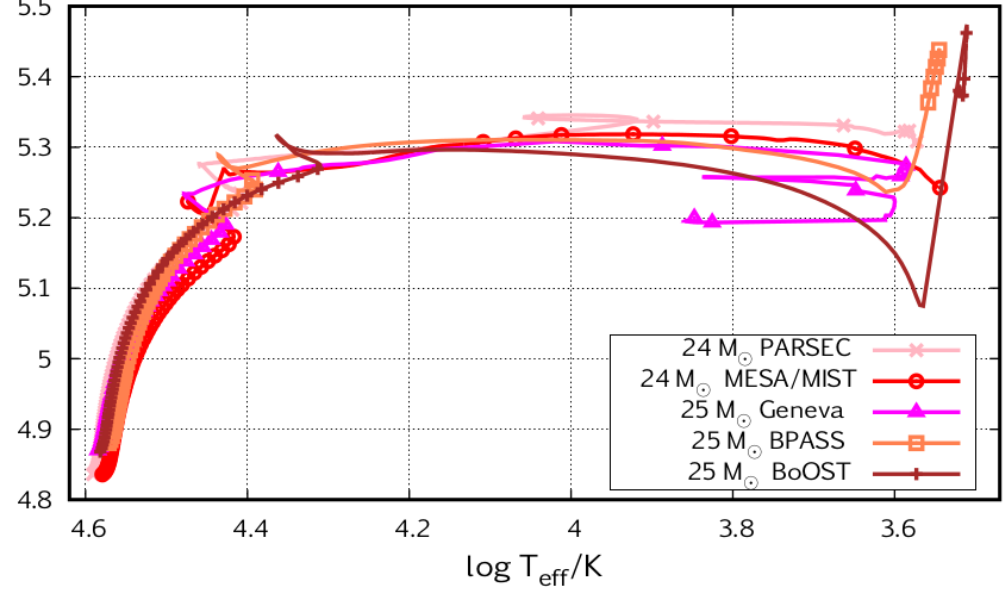
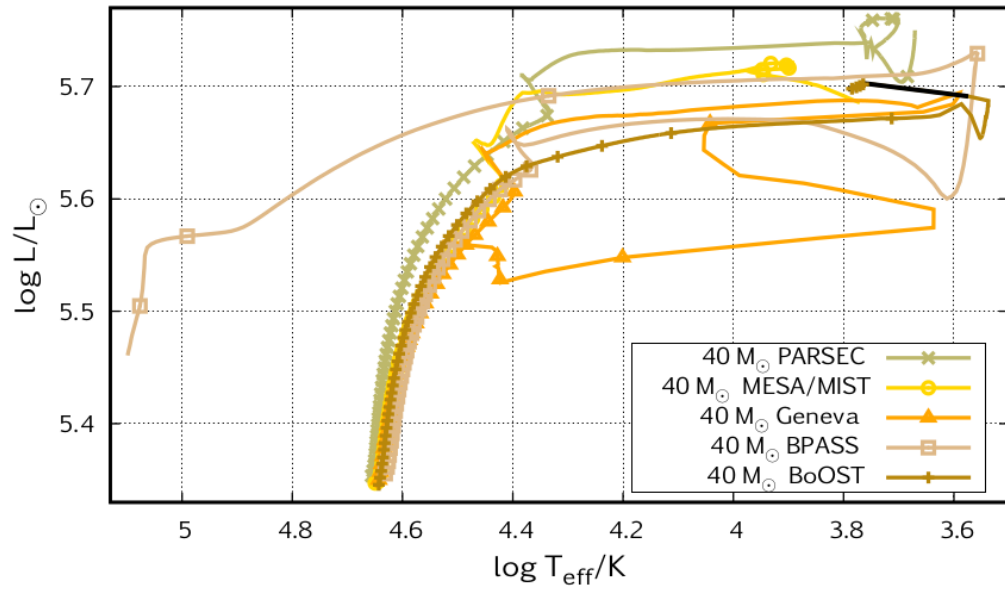
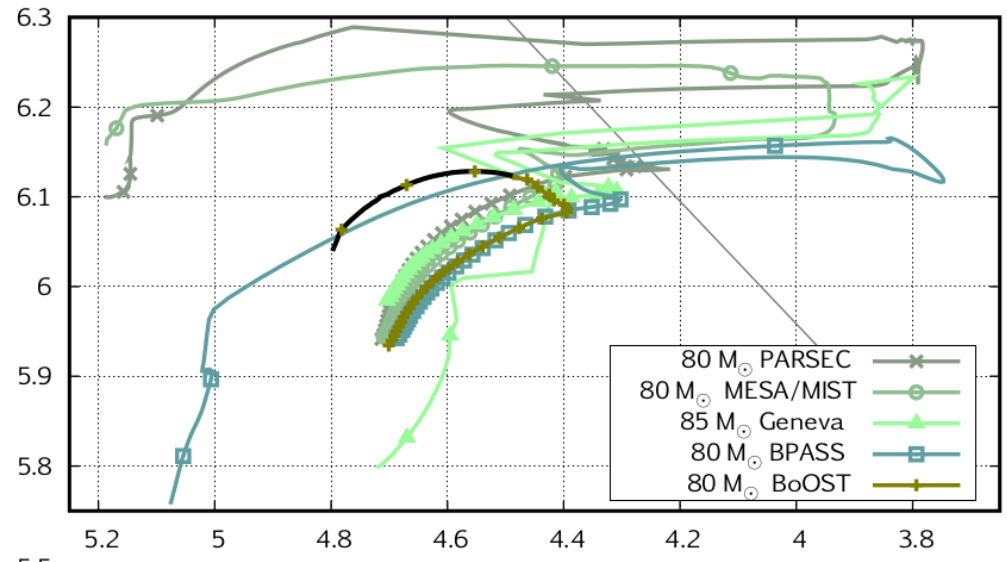
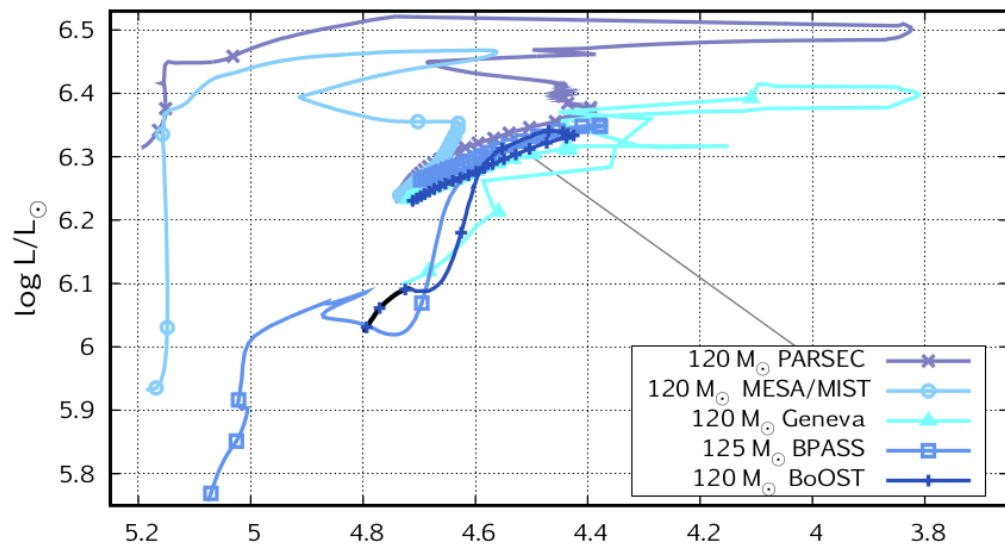
CORE

ENVELOPE

# How do the codes deal with that?

- several “tricks” in the literature
  - various codes use various tricks & methods
  - cf. Agrawal (*PhD Thesis*), Agrawal & Szécsi+22 (MNRAS)
- PARSEC (‘Padova’) artificially limiting the temp. gradient
- MIST (MESA) MLT++ formalism (*limiting the superadiabacity\**)  
=changing how convection\*\* is treated \*difference between the isothermal and adiabatic temperature gradient
- ‘Geneva’ } artificially enhanced mass loss at the right moment
- BPASS }
- BoOST (‘Bonn’) inflated envelope & post-processing with ‘normal’ mass loss





P. Agrawal (2021, *PhD thesis*)  
 Agrawal & Szécsi et al. (2022, MNRAS)

# Ionizing flux...

**Table 2.** Time averaged ionizing photon number flux [ $\text{s}^{-1}$ ] in the Lyman continuum emitted by the stellar models during their lives *on average*, cf. Section 4.2. The last column provides the amount of Lyman radiation (number of photons [ $\text{s}^{-1}$ ]) that a  $10^7 M_{\odot}$  population (e.g. a starburst galaxy or a young massive cluster in the Milky Way) containing these massive stars would emit.

$M_{\text{ini}} [M_{\odot}]$	24/25	40	80/85	120/125	pop.
PARSEC	$3.7 \times 10^{48}$	$1.3 \times 10^{49}$	$5.5 \times 10^{49}$	$1.0 \times 10^{50}$	$1.08 \times 10^{54}$
MIST	$3.3 \times 10^{48}$	$1.5 \times 10^{49}$	$5.1 \times 10^{49}$	$1.1 \times 10^{50}$	$1.06 \times 10^{54}$
Geneva	$3.5 \times 10^{48}$	$1.2 \times 10^{49}$	$5.1 \times 10^{49}$	$8.5 \times 10^{49}$	$9.90 \times 10^{53}$
BPASS	$3.6 \times 10^{48}$	$1.3 \times 10^{49}$	$4.5 \times 10^{49}$	$7.7 \times 10^{49}$	$9.34 \times 10^{53}$
BoOST	$3.7 \times 10^{48}$	$1.2 \times 10^{49}$	$4.2 \times 10^{49}$	$6.9 \times 10^{49}$	$8.89 \times 10^{53}$

up to 18% difference!

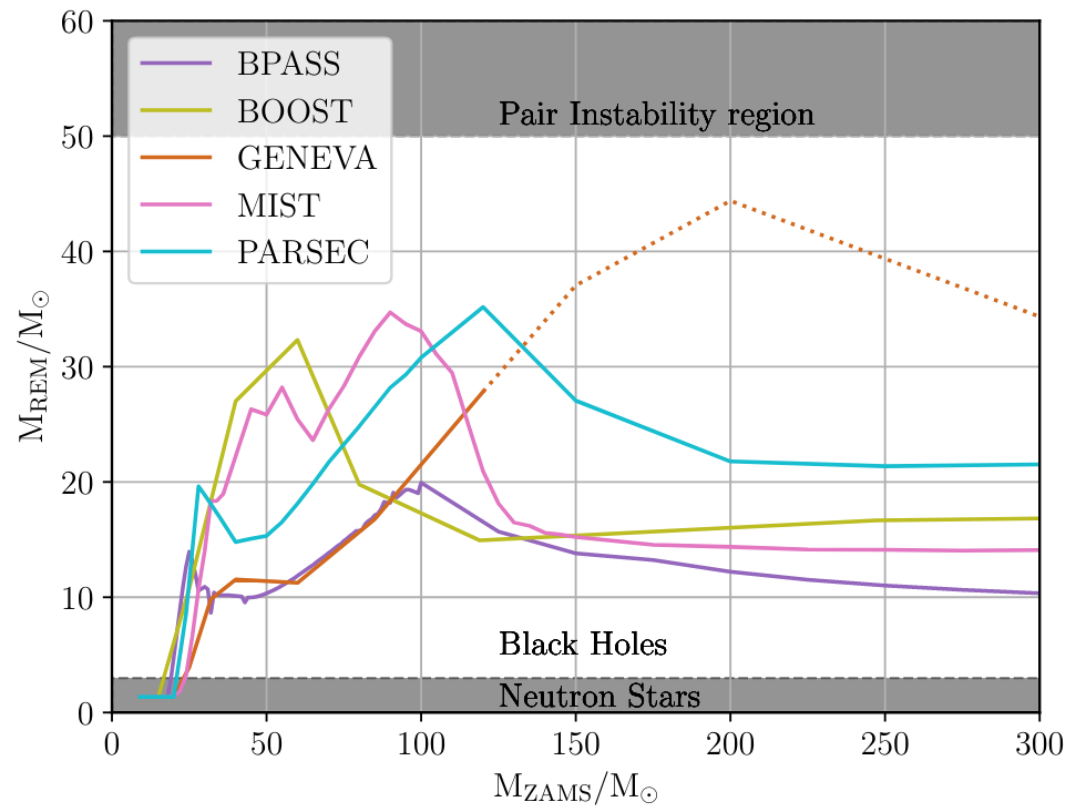
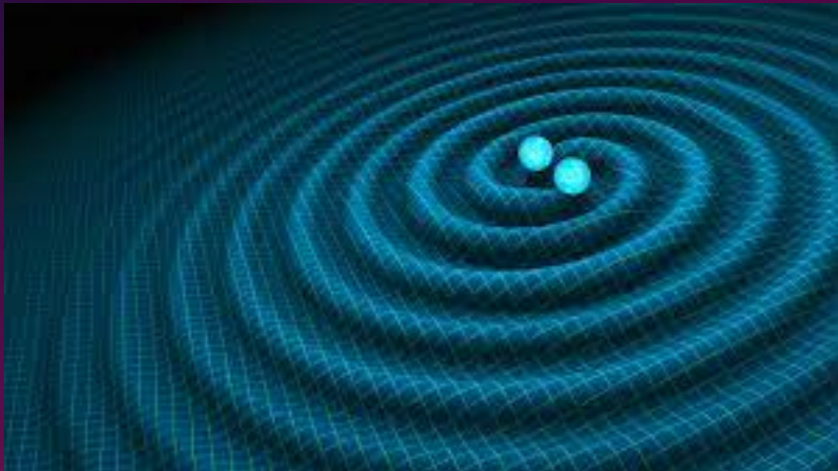


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# Remnant mass...

## Gravitational waves: compact object mergers (e.g. black holes)



**Figure 2.** Mass of stellar remnant as a function of the initial mass of the star (near-solar composition). Differences in the assumptions in massive star modelling can cause a variation of up to  $20 M_{\odot}$  in the remnant masses between simulations. Choosing to apply one of these simulations over the others in e.g. gravitational-wave event rate predictions can lead to strikingly different results.

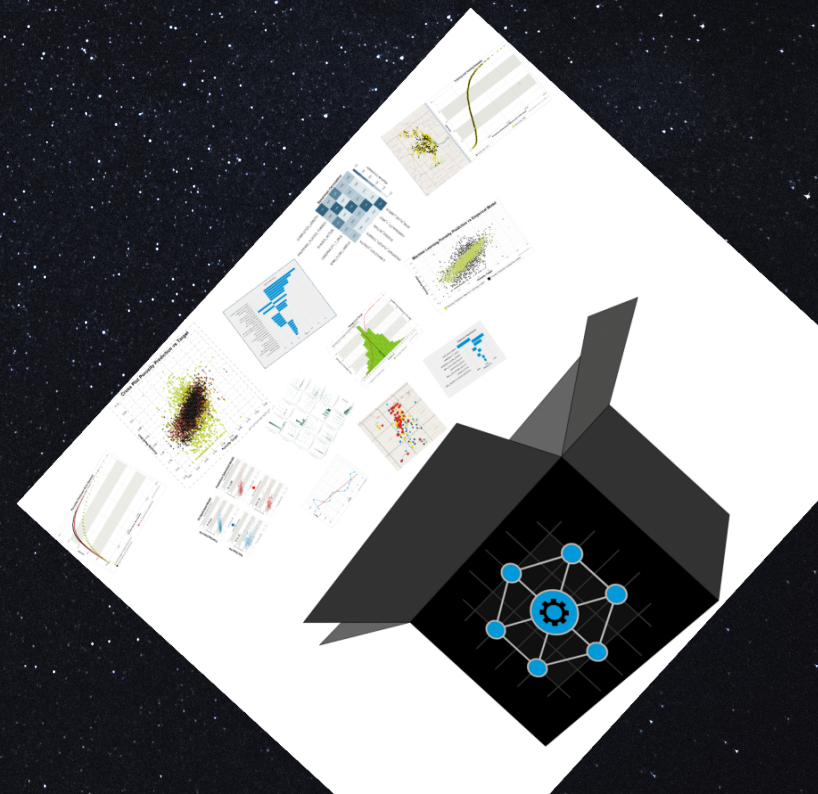
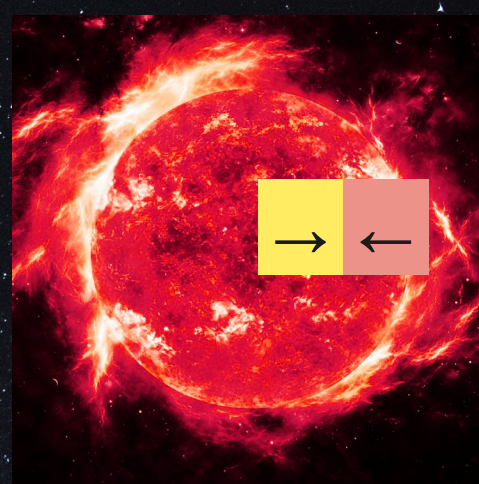
**up to  $20 M_{\odot}$  difference!**

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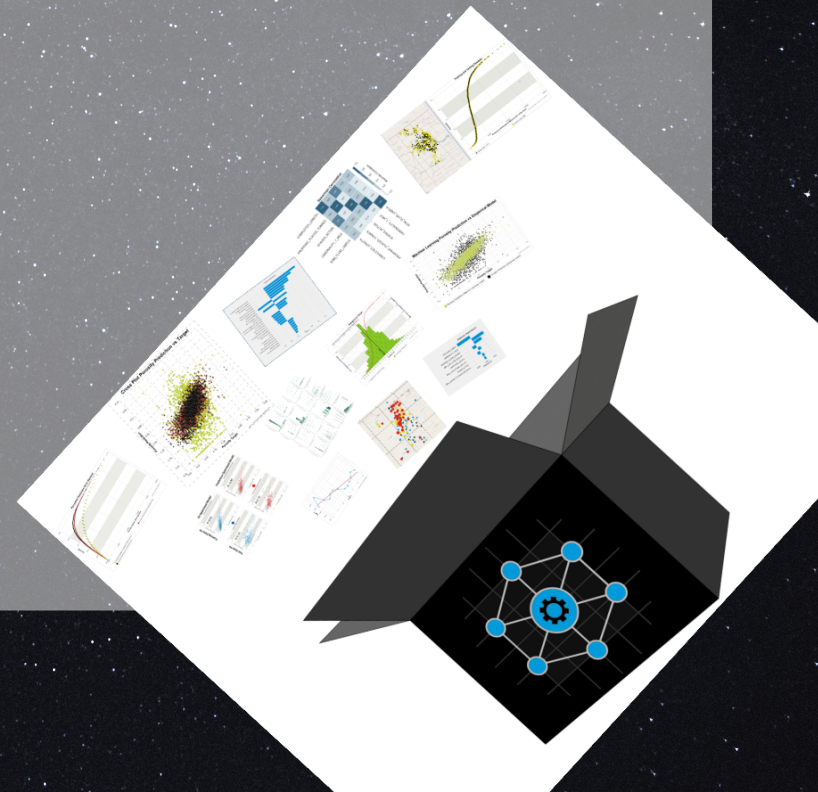
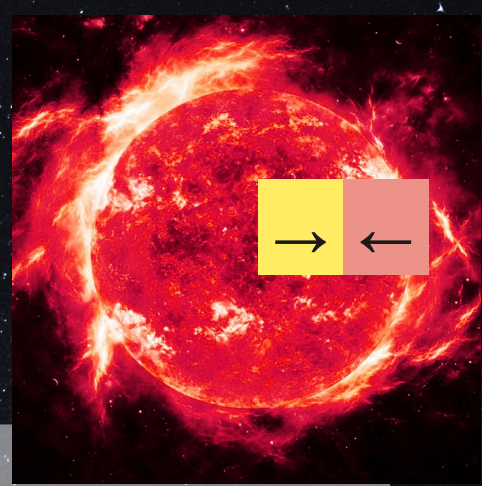


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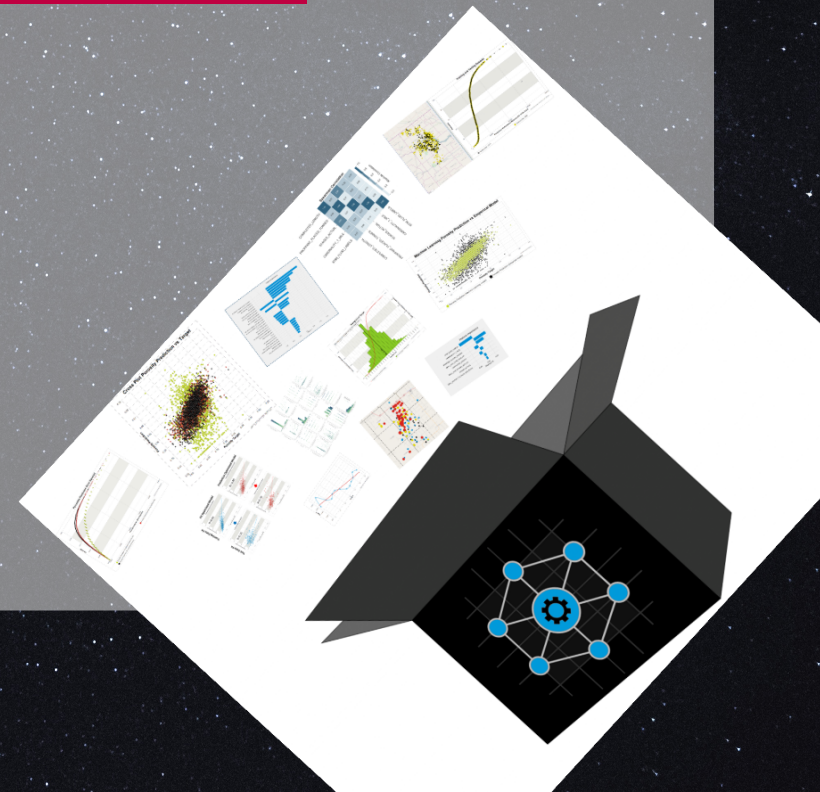
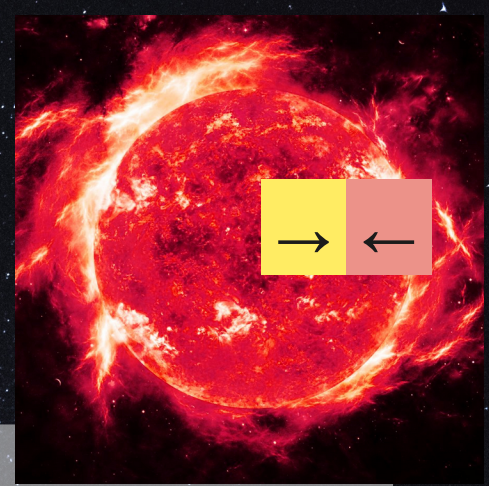
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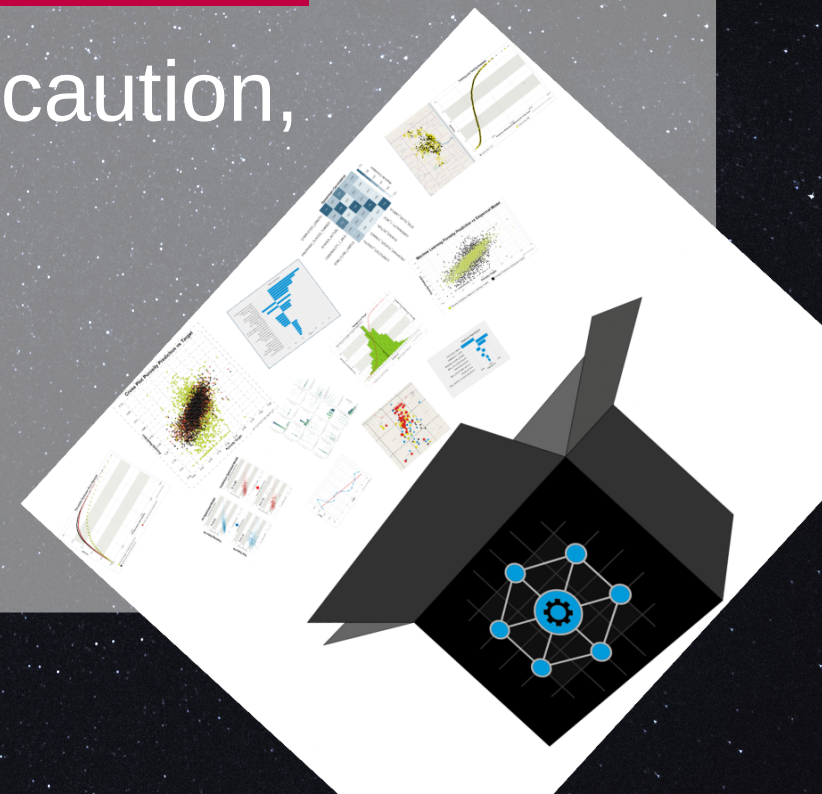
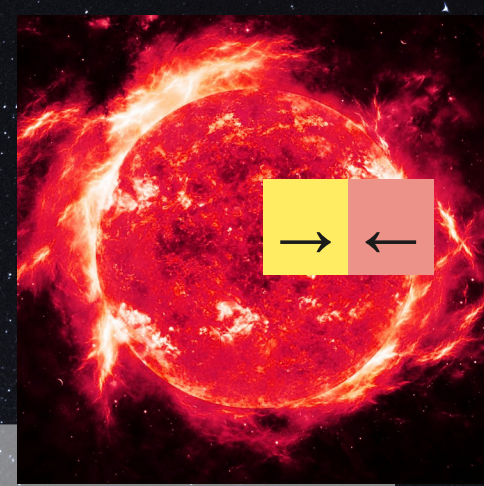
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**Thanks!**

